Rogue wave and a pair of resonance stripe solitons to a reduced generalized (3+1)-dimensional KP equation

Xiaoen Zhang, Yong Chen*, Xiaoyan Tang

Shanghai Key Laboratory of Trustworthy Computing, East China Normal University, Shanghai, 200062, China

Abstract

Based on the bilinear operator and symbol calculation, some lump solutions are presented, rationally localized in all directions in the space, to a reduced (3+1)dimensional KP equation. The lump solutions all contain six parameters, four of which must cater to the non-zero conditions so as to insure the analyticity and rational localization, while the others are free. Then the interaction between lump soliton and one stripe soliton is described and the result shows that the lump soliton will be drowned or swallowed by the stripe soliton. Furthermore, we extend this method to a new combination of positive quadratic function and hyperbolic functions. Especially, it is interesting that a rogue wave is found to be aroused by the interaction between lump soliton and a pair of resonance stripe solitons. By choosing the values of the parameters, the dynamic properties of lump solution, interaction between lump soliton and one stripe soliton, rogue wave, generated by the interaction between lump soliton and a pair of resonance stripe solution, interaction between lump soliton and a pair of resonance stripe solution, interaction between lump soliton and a pair of resonance solitons, are shown graphically.

Keywords: lump solution, interaction, a pair of resonance stripe solitons, rogue wave

1. Introduction

In soliton theory, the study to the integrability of nonlinear equation is always a hot topic. As to the integrable equations, there are many methods to study their solutions, such as the classic inverse scattering method[1], Bäcklund transformation[2, 3], Darboux transformation[4], Hirota bilinear methods[5, 6, 7], and variable separation approaches[8, 9, 10]. Among these methods, the Hirota bilinear method is widely popular due to its simplicity and directness. Recently, rogue wave solution[11](being as a special solution of the rational solution) attracts a lot of attention, which was first used to describe the momentous disastrous ocean waves. Its lethality is very

strong and can lead to devastating impact to the navigation. There are many ways to get the rogue wave solution, such as generalized Darboux transformation[12], bilinear method[13, 14], and so on. In contract to the rogue wave solution, lump solution is a special kind of rational solution, rationally localized in all directions in the space. In 2002, Lou et.al studied the lump solution with the variable separation method[15]. Recently, Ma proposed the positive quadratic function to get the lump solution. Special examples of lump solutions have been found, such as the KPI equation[16, 17, 18], BKP equation[19], the p-gKP and p-gBKP equations[20], Boussinesq equation[21] and so on.

More importantly, it will happen collision among different solitons. There are two kinds of collision, either elastic or inelastic. It is reported that lump solutions will keep their shapes, amplitudes, velocities after the collision with soliton solutions, which means the collision is completely elastic[22]. While many other collisions are completely inelastic, for instance, Becker et.al studied the inelastic collision of solitary waves in anisotropic Bose-Einstein condensates[23], Tan discussed the rational breather wave swallowed by kink wave[24], Tang showed the lump solution drowned by a stripe solution[25]. On the basis of different conditions, the collision will change essentially.

The main purpose of this paper is to study the lump solution, the interaction of lump soliton and one stripe soliton to a generalized (3+1)-dimensional KP equation[26]

$$(u_t + h_1 u u_x + h_2 u_{xxx} + h_3 u_x)_x + h_4 u_{yy} + h_5 u_{zz} = 0,$$
(1)

when $h_1 = -1$, $h_2 = -\frac{1}{3}$, $h_3 = 1$, $h_4 = 1$, $h_5 = -\frac{2}{3}$. Moreover, we extend this method for a combination of positive quadratic function and hyperbolic cosine, then it appears a strange phenomenon, we recall it rogue wave(also be called ghost soliton)from its mathematics expression, generated by the interaction of lump soliton and a pair of resonance solitons in a evolutionary process.

The structure of this paper is as follows: Based on the bilinear operator, Sec.2 obtains the lump solution with the positive quadratic function. Sec.3 list the lump soliton, one stripe soliton and their interaction by using method of the collection of positive quadratic function and exponential function. In the end, we extend the positive quadratic function to a combination of hyperbolic cosine function, during the evolutionary process, there exists a rogue wave as the motivation of the interaction between lump soliton and a pair of resonance solitons.

2. Lump solution of a reduced generalized (3+1)-dimensional KP equation

When $h_1 = -1, h_2 = -\frac{1}{3}, h_3 = 1, h_4 = 1, h_5 = -\frac{2}{3}$, Eq.(1) becomes

$$(u_t - uu_x - \frac{1}{3}u_{xxx} + u_x)_x + u_{yy} - \frac{2}{3}u_{zz} = 0.$$
 (2)

With the transformation $u = 4(\ln f)_{xx}$, its bilinear equation can be presented

$$(D_x D_t - \frac{1}{3}D_x^4 + D_x^2 + D_y^2 - \frac{2}{3}D_z^2)f \cdot f = 0,$$
(3)

When z = x, Eq. (3) becomes the following formula

$$(D_x D_t - \frac{1}{3}D_x^4 + \frac{1}{3}D_x^2 + D_y^2)f \cdot f$$

$$= 2f_{xt}f - 2f_xf_t - \frac{2}{3}f_{xxxx}f + \frac{8}{3}f_{xxx}f_x - 2f_{xx}^2 + \frac{2}{3}f_{xx}f - \frac{2}{3}f_x^2 + 2f_{yy}f - 2f_y^2,$$
(4)

which can be changed KP equation with some variable transformations.

Assume

$$f = g^{2} + h^{2} + a_{9}, g = a_{1}x + a_{2}y + a_{3}t + a_{4}, h = a_{5}x + a_{6}y + a_{7}t + a_{8},$$
(5)

where $a_i, 1 \le i \le 9$ are parameters to be determined. By substituting f into Eq.(3), with a direct calculation, these parameters can be expressed:

$$\begin{cases} a_3 = \frac{-a_1^3 + (3a_6^2 - a_5^2 - 3a_2^2)a_1 - 6a_5a_6a_2}{3a_5^2 + 3a_1^2}, a_7 = \frac{-a_5^3 + (3a_2^2 - a_1^2 - 3a_6^2)a_5 - 6a_1a_6a_2}{3a_5^2 + 3a_1^2} \\ a_9 = \frac{(a_1^2 + a_5^2)^3}{(a_1a_6 - a_2a_5)^2}, \end{cases}$$

which should be satisfied

$$a_1 a_5 \neq 0$$
, and $a_1 a_6 - a_2 a_5 \neq 0$, (7)

in order to insure the analytical and positive of f, meanwhile,

$$\begin{vmatrix} a_1 & a_2 \\ a_5 & a_6 \end{vmatrix} \neq 0 \tag{8}$$

guarantee the rationally localization of u in all directions in the (x, y)-plane.

In return, substitute Eq.(6) into Eq.(5) and generate a class of positive, analytical f:

$$f = (a_1x + a_2y + \frac{-a_1^3 + (3a_6^2 - a_5^2 - 3a_2^2)a_1 - 6a_5a_6a_2}{3a_5^2 + 3a_1^2}t + a_4)^2 + \frac{(a_1^2 + a_5^2)^3}{(a_1a_6 - a_2a_5)^2} + (a_5x + a_6y + \frac{-a_5^3 + (3a_2^2 - a_1^2 - 3a_6^2)a_5 - 6a_1a_6a_2}{3a_5^2 + 3a_1^2}t + a_8)^2,$$
(9)

then the solution of u can be written through the transformation $u = 4(\ln f)_{xx}$

$$u = \frac{8(a_1^2 + a_5^2)}{(a_1x + a_2y + a_3t + a_4)^2 + (a_5x + a_6y + a_7t + a_8)^2 + a_9} - \frac{16((a_1x + a_2y + a_3t + a_4)a_1 + (a_5x + a_6y + a_7t + a_8)a_5)^2}{((a_1x + a_2y + a_3t + a_4)^2 + (a_5x + a_6y + a_7t + a_8)^2 + a_9)^2}.$$
(10)

In this class of lump solution, Eq.(7) and Eq.(8) should satisfy, its expicitly pictures are showed when y = 0 and t = 0 respectively. Moreover, its density plots are given and the shape of lump solution is clearer.

3. The interaction between lump soliton and one stripe soliton

We want to study the collision between the lump soliton and one stripe soliton. In sec.2, based on the positive quadratic function, its lump solution are presented. In this section, make f as a combination of positive quadratic function and one exponential function, that is

$$f_1 = m_1^2 + n_1^2 + l_1 + a_9, (11)$$

where

$$m_1 = a_1x + a_2y + a_3t + a_4, n_1 = a_5x + a_6y + a_7t + a_8, l_1 = ke^{k_1x + k_2y + k_3t},$$

through substituting Eq.(11) into Eq.(4) and symbol calculation, these parameters are calculated:

$$a_{1} = -\frac{a_{6}}{k_{1}}, a_{2} = 0, a_{3} = -\frac{a_{6}(-1+3k_{1}^{2})}{3k_{1}}, a_{5} = 0, a_{7} = 0, a_{9} = \frac{a_{6}^{2}}{k_{1}^{4}}, k_{2} = 0, k_{3} = \frac{1}{3}(k_{1}^{3}-k_{1}),$$
(12)

which needs to satisfy conditions

$$k_1 \neq 0, \quad k > 0, \tag{13}$$

to make the corresponding solutions f is positive, analytical and guarantee the localization of u in all directions in the (x, y)-plane.

Based on the transformation $u = 4(\ln f_1)_{xx}$, the solution of Eq.(4) will be got again

$$u = \frac{4(2a_1^2 + 2a_5^2 + k_1^2l)}{f_1} - \frac{4(2a_1m + 2a_5n + k_1l)^2}{f_1^2},$$
(14)

where

$$f_1 = \left(-\frac{a_6x}{k_1} - \frac{a_6(-1+3k_1^2)t}{3k_1} + a_4\right)^2 + \left(a_6y + a_8\right)^2 + ke^{k_1x + \left(\frac{k_1^3}{3} - \frac{k_1}{3}\right)t} + \frac{a_6^2}{k_1^4},$$
$$m_1 = -\frac{a_6x}{k_1} - \frac{a_6(-1+3k_1^2)t}{3k_1} + a_4, n_1 = a_6y + a_8, l_1 = ke^{k_1x + \left(\frac{k_1^3}{3} - \frac{k_1}{3}\right)t}.$$

By choosing appropriate values of these parameters, the dynamic graphs of collision between the lump soliton and one stripe soliton are showed in Fig. ??, Fig. ??: Whereafter, its density plots are shown in Fig. ?? Fig. ?? (a) shows there is one lump soliton and one stripe soliton, the energy of lump soliton is stronger than stripe soliton, when t is up to 0, lump soliton begins to be swallowed by stripe soliton step by step, its energy begin to transfer into the stripe soliton gradually, until it is swallowed by the stripe soliton completely, these two kinds of solitons roll into one soliton and continue to spread.

4. Rogue wave and a pair of resonance solitons

Based on the collision of lump soliton and one stripe soliton, we begin to discuss the collision of lump soliton and two stripe solitons. Take f as the combination of positive quadratic function and two exponential functions in virtue of method to seek N-soliton solutions of bilinear form, that is

$$f_2 = m_2^2 + n_2^2 + kg_2 + k_4h_2 + k_8g_2h_2, (15)$$

where

$$m_2 = a_1 x + a_2 y + a_3 t + a_4, n_2 = a_5 x + a_6 y + a_7 t + a_8, g_2 = e^{k_1 x + k_2 y + k_3 t}, h_2 = e^{k_5 x + k_6 y + k_7 t},$$

by substituting Eq.(15) into Eq.(4), we can get

$$\begin{cases} a_1 = \frac{a_6}{k_5}, a_2 = 0, a_3 = \frac{a_6(-1+3k_5^2)}{k_5}, a_5 = 0, a_7 = 0, a_9 = \frac{a_6^2}{k_5^4}, k = \frac{a_6^2k_8}{k_4k_5^4}, \quad (16)\\ k_1 = -k_5, k_2 = 0, k_3 = \frac{k_1^4 - 3k_2^2 - k_1^2}{3k_1}, k_6 = 0, k_7 = \frac{k_5^4 - k_5^2 - 3k_6^2}{3k_5}, \end{cases}$$

the results indicates these two exponential functions are a pair of resonance solitons, by some transformations, it can be changed into a hyperbolic cosine function, hence, reinstall f as the following formula:

$$f_3 = m_3^2 + n_3^2 + k \cosh(k_1 x + k_2 y + k_3 t) + a_9, \tag{17}$$

where

$$m_3 = a_1 x + a_2 y + a_3 t + a_4, n_3 = a_5 x + a_6 y + a_7 t + a_8,$$

once again, substitute Eq.(17) into Eq.(4), with a complex symbol calculation, the relations of these parameters are

$$\begin{cases} a_1 = \frac{a_6}{k_1}, a_3 = \frac{a_6^2(-1+3k_1^2) - 3a_2^2k_1^2}{3k_1a_6}, a_5 = 0, a_7 = -2a_2k_1, a_9 = \frac{k^2k_1^8 + 4a_6^4}{4a_6^2k_1^4} (18)\\ k_2 = \frac{a_2k_1^2}{a_6}, k_3 = \frac{(k_1^3 - k_1)a_6^2 - 3a_2^2k_1^3}{3a_6^2}, \end{cases}$$

which needs to satisfy

$$k_1 \neq 0, \quad a_6 \neq 0, \quad k > 0,$$
 (19)

to guarantee the corresponding solutions f_3 is positive, analytical and insure the localization of u in all directions in the (x, y)-plane.

Still, institute Eq.(18) into Eq.(17), with the transformation $u = 4(\ln f_3)_{xx}$, we can obtain the solution of u

$$u = \frac{4(2a_1^2 + 2a_5^2 + k\cosh(k_1x + k_2y + k_3t)k_1^2)}{f_3} - 4\frac{(2a_1m_3 + 2a_5n_3 + kk_1\sinh(k_1x + k_2y + k_3t))^2}{f_3^2},$$
(20)

where

$$\begin{cases} f_3 = \left(\frac{a_6}{k_1}x + a_2y + \frac{a_6^2(-1+3k_1^2) - 3a_2^2k_1^2}{k_1a_6}t + a_4\right)^2 + \left(a_6y - 2a_2k_1t + a_8\right)^2 + \frac{k^2k_1^8 + 4a_6^4}{4a_6^2k_1^4} \\ +k\cosh(k_1x + \frac{a_2k_1^2}{a_6}y + \frac{(k_1^3 - k_1)a_6^2 - 3a_2^2k_1^3}{3a_6^2}t), \quad n_3 = a_6y - 2a_2k_1t + a_8, \\ m_3 = \frac{a_6}{k_1}x + a_2y + \frac{a_6^2(-1+3k_1^2) - 3a_2^2k_1^2}{k_1a_6}t + a_4. \end{cases}$$
(21)

According to the expression of f_3 , n_3 , m_3 , asymptotic property of lump solution and a pair of resonance solitons are analyzed.

Take

$$\xi_1 = \frac{a_6}{k_1}x + a_2y + \frac{a_6^2(-1+3k_1^2) - 3a_2^2k_1^2}{k_1a_6}t + a_4, \quad \xi_2 = a_6y - 2a_2k_1t + a_8$$

$$\xi_3 = k_1x + \frac{a_2k_1^2}{a_6}y + \frac{(k_1^3 - k_1)a_6^2 - 3a_2^2k_1^3}{3a_6^2}t,$$

with a comparison for ξ_1, ξ_2, ξ_3 , it is proved

$$\xi_1 = \xi_3 \frac{a_6}{k_1^2} + \frac{2a_6k_1t}{3} + a_4, \lim_{t=\pm\infty} \frac{\xi_1^2}{\xi_2^2} = \frac{(3a_2^2k_1^2 - 3a_6^2k_1^2 + a_6^2)^2}{k_1^4a_6^2a_2^2}$$

due to $t = \pm \infty$, ξ_1, ξ_2 are same order, we only need to compare ξ_1^2 , $\cosh \xi_3$, if suppose ξ_3 is a constant, then ξ_1 is a combination of scale change and time displacement for ξ_3 . When $t = \pm \infty$,

$$\lim_{t=\pm\infty} \frac{\xi_1^2}{\cosh(\xi_3)} = \lim_{t=\pm\infty} \frac{(\xi_3 \frac{a_6}{k_1^2} + \frac{2a_6k_1t}{3} + a_4)^2}{\cosh(\xi_3)} = 0.$$

so we can come to a conclusion, when $t = \pm \infty$, there only are a pair of resonance solitons, when t becomes little, lump soliton property is more obvious, which can be seen in Fig. ??. Fig. ?? (a) indicates there are a pair of resonance solitons, lump soliton is in a invisible place, similar to a ghoston, (b) shows when t = -2, lump soliton appears gradually, it attaches to one of the resonance stripe soliton. More importantly, because of energy conservation, the shapes of these two resonance solitons change at the same time and in the same location, one position appears a lump soliton, the other relative position appears a sunk envelope. When t = 0, there exists a rogue wave, derived from lump soliton, is located in the middle of these two resonance solitons and link them with each other, then its carrier begin to transfer, until it attaches to other stripe soliton successfully and out of our vision. Fig. ?? (a) shows there are only a pair of resonance stripe solitons, lump soliton maybe as a ghoston, which is hidden in one of the stripe solitons, (b) shows lump soliton appears in one of stripe soliton, when t is up to 0, lump soliton's energy reaches up to the maximum, which presents the property of rogue wave. Whereafter, its energy transfer into the other stripe until disappearing. This whole progress can be regarded as a appearing of the rogue wave. Then its sectional drawing and vertical view are showed in Fig. ?? Fig. ?? respectively green line represents t=-10, blue line represents t=10, red line represents t=0. Obviously, the amplitude of t=0 is about five times than $t=\pm 10$ as well as its appearing time is short, which is cater to the properties of rogue wave. So this kind of lump solution can also be called rogue wave, which is aroused by the interaction between lump soliton and a pair of resonance solitons.

5. Discussion

Based on the Hirota formulation and symbol calculation, we research a reduced (3+1)-dimensional KP equation (2). First, its lump soliton is got and the analyticity, localization of the resulting is guaranteed by the non-zero determinant and some constraints to the parameters. But not every Eq.(1) have lump soliton, when $h_1 = 1, h_2 = 1, h_3 = 1, h_4 = 1, h_5 = 1$, the bilinear formulation of Eq.(1) is

$$(D_x D_t + D_x^4 + D_x^2 + D_y^2 + D_z^2)f \cdot f = 0.$$
(22)

Once more, set z = x and assume f as (5), then the quadratic function of f is

$$f = (a_1x + a_2y + \frac{-a_1^3 + (-2a_5^2 - a_2^2 + a_6^2)a_1 - 2a_5a_6a_2}{a_1^2 + a_5^2} + a_4)^2 + (-\frac{3(a_1^2 + a_5^2)^3}{(a_1a_6 - a_2a_5)^2}) + (a_5x + a_6y + \frac{-a_5^3 + (-2a_1^2 - a_6^2 + a_2^2)a_5 - 2a_1a_6a_2}{a_1^2 + a_5^2})^2,$$

we can see $a_9 < 0$ in the expression of f, so there doesn't exist lump soliton, let alone the collision with other solitons.

Second, we describe the collision in lump soliton and one stripe soliton, the dynamic property is presented in Fig. ??, in the beginning, lump soliton keep its types, energy, spread with a steady rate, but when it meets the stripe soliton, it will interact with the stripe soliton, but lump soliton is swallowed in the end.

Furthermore, we characterize the rogue wave and a pair of resonance solitons, their dynamic property is showed in Fig. ??, in the beginning, there have a pair of resonance stripe solitons, but their amplitudes are different, we can see there should exist a mysterious soliton. Then we discover there appears a small lump soliton which attaches to one of stripe soliton, and it sinks a wave packet in the homologous location for the other stripe soliton. As times goes on, the interaction between this lump soliton and resonance solitons becomes more vigorous, in a special time, its amplitude is up to the maximum, which can be seen in Fig. ??. Its amplitude changes greatly and its occurrence time is short, which caters to the character of rogue wave, so this progress can be called the generation of the rogue wave. Whereafter, rogue wave degenerate into a lump soliton and attaches to the resonance soliton again until disappearing. As we all know, rogue wave was first be found to the KP type equation.

In addition, we should try to discuss the interactions between lump soliton and other kinds of solitons. Moreover, we can expand this method to discrete equations and study their lump solution. These problems will be more interesting and be worthy of discussing.

Acknowledgment

We would like to express our sincere thanks to S Y Lou, W X Ma and other members of our discussion group for their valuable comments. The project is supported by the Global Change Research Program of China(No.2015CB953904), National Natural Science Foundation of China (No. 11275072, 11435005, 11675054, 11275123 and 11675055). The Network Information Physics Calculation of basic research innovation research group of China (No.61321064), and Shanghai Collaborative Innovation Center of Trustworthy Software for Internet of Things (No.ZF1213).

References

- M. J. Ablowitz, P. A. Clarkson, Solitons, Nonlinear Evolution Equations and Inverse Scattering, Cambridge University Press, New York, 1991.
- [2] C. Rogers, W. F. Shadwick, Bäcklund Transformations and Their Applications, Academic Press, London, 1982.
- [3] R. M. Miura, Bäcklund Transformation, Springer-Verlag, Berlin, 1978.
- [4] V.B. Matveev, M.A. Salle, Darboux Transformations and Solitons, Springer-Verlag, Berlin, 1991.
- [5] R. Hirota, The direct method in soliton theory, Cambridge University Press, 2004.
- [6] R. Hirota, Exact soliton of the Korteweg-de Vries equation for multiple collisions of solitons, Physical Review Letters, Vol.27(1971): 1192-1194.
- [7] R. Hirota, Exact N-soliton solutions of the wave equation of long waves in shallow water and in nonlinear lattices, Journal of Mathematical Physics, Vol.14(1973) 810-814.
- [8] S. Y. Lou, C. L. Chen, X. Y. Tang, (2+1)-dimensional (M+N) -component AKNS system: Painlev integrability, infinitely many symmetries, similarity reductions and exact solutions. Journal of Mathematical Physics, Vol.43(2002): 4078-4108.
- [9] X. Y. Tang, S. Y. Lou, Y. Zhang, Localized excitations in (2+1)-dimensional systems, Physical Review E, Vol.66(2001): 046601.

- [10] X. Y. Tang, C. L. Chen, S. Y. Lou, Localized solutions with chaotic and fractal behaviours in a (2+1)-dimensional dispersive long-wave system, Journal of Physics A: Mathematical and General, Vol.35(2002): L293-L301.
- [11] C. Garrett, J. Gemmrich, Rogue waves, Physics today, Vol.62(2009): 62-63.
- [12] X. Wang, Y. Q. Li, F. Huang, Y. Chen, Rogue wave solitons of AB system, Communications in Nonlinear Science and Numerical Simulation, Vol.20(2015): 434-442.
- [13] J. C. Chen, Y. Chen, B. F. Feng, K. I. Maruno, Multi-dark soliton solutions of the twodimensional multi-component yajima-oikawa systems, Journal of the Physical Society of Japan, 84(2015): 034002.
- [14] J. C. Chen, Y. Chen, B. F. Feng, K. I. Maruno, Rational solutions to twoand one-dimensional multicomponent yajima-oikawa systems, Physics Letters A, 379(2015): 1510-1519.
- [15] S. Y. Lou, X. Y. Tang, Nonlinear mathematical physics method, Academic Press, Beijing, 2006.
- [16] W. X. Ma, Lump solutions to the Kadomtsev-Petviashvili equation, Physics Letters A, Vol.379(2015): 1975-1978.
- [17] S.V. Manakov, V.E. Zakharov, L.A. Bordag, A.R. Its and V.B. Matveev, Twodimensional solitons of the Kadomtsev-Petviashvili equation and their interaction, Physics Letters A, Vol.63(1977): 205-206.
- [18] R.S. Johnson and S. Thompson, A solution of the inverse scattering problem for the Kadomtsev-Petviashvili equation by the method of separation of variables, Physics Letters A, Vol.66(1978): 278-281.
- [19] J. Y. Yang, W. X. Ma, Lump solutions to the BKP equation by symbolic computation, International Journal of Modern Physics B, Vol.30(2016): 1640028.
- [20] W. X. Ma, Z. Y. Qin and X. Lü, Lump solutions to dimensionally reduced p-gKP and p-gBKP equations, Nonlinear Dynamics, Vol.84(2016): 923-931.
- [21] H. C. Ma, A. P. Deng, lump solution of (2+1)-dimensional Boussinesq equation, Communications in Theoretical Physics, Vol.65(2016): 546-552.
- [22] A. S. Fokas, D. E. Pelinovsky, C. Sulem, Interaction of lumps with a line soliton for the DSII equation, Physica D, Vol.152-153(2001): 189-198.

- [23] C. Becker, K. Sengstock, P. Schmelcher, P. G. Kevrekidis, R. C. González, Inelastic collisions of solitary waves in anisotropic Bose-Einstein condensates: sling-shot events and expanding collision bubbles, New Journal of Physics, Vol.15(2013): 113028.
- [24] W. Tan, Z. D. Dai, Dynamics of kinky wave for (3+1)-dimensional potential Yu-Toda-Sasa-Fukuyama equation, Nonlinear Dynamics, Vol.85(2016): 817-823.
- [25] Y. N. Tang, S. Q. Tao, Q. Guan, Lump solitons and the interaction phenomena of them for two classes of nonlinear evolution equations, Computers and Mathematics with Applications, (2016).
- [26] J. M. Tu, S. F. Tian, M. J. Xu, X. Q. Song, T. T. Zhang, Bäklund transformation, infinite conservation laws and periodic wave solutions of a generalized (3+1)-dimensionalnonlinear wave in liquid with gas bubbles, Nonlinear Dynamics, Vol.83(2016): 1199-1215.
- [27] W. X. Ma, T. C. Xia, Pfaffianized systems for a generalized Kadomtsev-Petviashvili equation, Physica Scripta, Vol.87(2013): 458-465.